

1. We use eq. from BSL I Sect 9.6 to get overall heat-transfer coef. U_o , but first we need h_o at inner surface.

$$Re = \frac{Dv\rho}{\mu} = \frac{0.6(2)1.1}{3.1 \cdot 10^{-5}} = 44,000$$

$$Pr = \frac{c_p \mu}{k} = \frac{1006(3.105)}{0.025} = 1.207$$

Re is large enough ($>20,000$) to use BSLK Eq. 14.3-16:

$$Nu = 0.026 Re^{0.8} Pr^{1/3} = 0.026(44000)^{0.8} (1.207)^{1/3} \quad \begin{array}{l} \uparrow \\ \text{see note} \\ \text{at end} \\ \text{of exam} \end{array}$$

$$= 143.5 = \frac{hD}{k} = \frac{h(0.6)}{0.025} \rightarrow h = 5.98$$

BSL I 9.6-31 (BSLK 10.3-31)

$$(\Gamma_o U_o)^{-1} = \left[\frac{1}{(0.3)(5.98)} + \frac{r_o(0.35/0.3)}{0.5} + \frac{1}{(0.35)(50)} \right] \quad (I)$$

$$= [0.557 + 0.308 + 0.057]$$

(resistance from heat transfer to inner wall) conduction thru wall heat transfer to air

$$U_o = \frac{1}{0.3} [0.922]^{-1} = 3.61 \text{ W/mK}$$

Eq. from class:

$$\frac{U_o D_o}{k} = r_o \left(\frac{T_f - T_{b1}}{T_f - T_{b2}} \right) Re Pr \frac{D_o}{4L}$$

$$\frac{(3.61)(0.6)}{0.025} = r_o \left(\frac{20 - 300}{20 - T_{b2}} \right) (44000)(1.207) \frac{0.6}{4 \times 50}$$

$$86.64 = r_o \left(\frac{20 - 300}{20 - T_{b2}} \right) 159$$

$$\frac{20 - 300}{20 - T_{b2}} = \exp\left(\frac{86.64}{159}\right) = 1.72$$

$$T_{b2} = 182^\circ \text{C}$$

b) Back to Eq. I above. Heat transfer to the air is an insignificant component in this problem; its resistance is small compared to the others. I could be off by 2 or 3 times with almost no effect. (In fact, I think I probably am wrong by more than that!)

Notes ρ and c_p for the smokestack don't matter for this problem.

2. Since T is assumed uniform, we can take a macro energy balance. Assume wire length L . (The value of L doesn't matter)

$$\text{generation: } \pi R^2 L S$$

$$\text{heat transfer at surface: } -(2\pi R L) A (T - T_0)^4$$

(> 0 because heat is lost)

$$\text{accumulation } \pi R^2 L \rho C_p \frac{dT}{dt}$$

$$(\pi R^2 L S) - (2\pi R L) A (T - T_0)^4 = \pi R^2 L \rho C_p \frac{dT}{dt}$$

$$RS - 2A(T - T_0)^4 = R\rho C_p \frac{dT}{dt}$$

$$\frac{dT}{dt} = \frac{1}{\rho C_p} \left(S - \frac{2A}{R} (T - T_0)^4 \right)$$

initial condition: $T = T_0$ at $t = 0$

3. This is unsteady conduction in a semi-infinite solid.

$$\text{BSLK Eq. 11.5-12: } q_x = \frac{K}{\sqrt{\pi t}} (T_i - T_0)$$

(where in our case conduction is in the x direction)

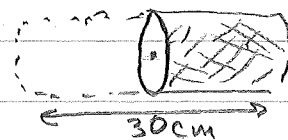
$$q = \frac{K}{\sqrt{\pi t}} \sqrt{\frac{K \rho C_p}{\pi t}} (T_i - T_0) = \sqrt{\frac{K \rho C_p}{\pi t}} (T_i - T_0)$$

$$-7000 = \sqrt{K} \left(\frac{11000 \cdot 130}{\pi \cdot 3600} \right)^{1/2} (20 - 100) = -900 \sqrt{K}$$

$$K = 60.5 \text{ W/m}^2\text{K}$$

4. a) This is a product of a slab + cylinder. The insulated surface makes it equivalent to a slab 30 cm thick.

The slowest spot to change is the center of the insulated surface, which is the center of the 30 cm slab.



$$b) \alpha = \frac{K}{\rho C_p} = \frac{23}{7820 \cdot 461} = 6.37 \cdot 10^{-6}$$

$$\alpha_{sl}: (\alpha t / R^2) = 6.37 \cdot 10^{-6} \times 60 / (0.05)^2 = 0.153$$

From BSLK Fig 11.5-2, at $r/R = 0$, $\frac{T_1 - T}{T_1 - T_0} = 0.64$

slab: $\alpha t / b^2 = 6.37 \cdot 10^{-6} (60) / (0.15)^2 = 0.02$

from Fig 11.5-1, no effect of insulated edge: $\frac{T_1 - T}{T_1 - T_0} = 1$

For solid, $\frac{T_1 - T}{T_1 - T_0} = 0.64 \cdot 1 = 0.64 = \frac{100 - T}{100 - 50} \rightarrow T = 68^\circ\text{C}$

- c) Internal conduction and heat transfer to the surface are in series. The slower process dominates. If the second calculation says the whole solid is at 90°C , then that process is ^{much} faster and therefore ^{it is} less important. The calculation in part (b) is more accurate.

Note on problem 1: in solving for h_p , one could also use

BSLK Fig 14.3-2 (instead of BSLK Eq. 14.3-16).

For $Re = 44,000$, one gets $0.0032 \approx \frac{h_{eq} D}{k} (Re)^{-1} (Pr)^{-1/3}$

$$0.0032 = \frac{h_{eq} (0.6)}{0.025} (44000)^{-1} (1.207)^{-1/3}$$

$$h_{eq} = 6.25 \rightarrow h_0 \text{ in eq. for } U_0.$$

Result is nearly the same as w/ Eq. 14.3-16.